

Rigid body time-stepping schemes in a quasi-static setting

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Abstract. We discuss how linear complementary problems (LCPs) can be used to simulate rigid-body systems in a quasi-static setting. LCP-based time-stepping schemes were successfully used in [1] in order to plan and control meso-scale manipulation tasks.

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1. Introduction

In [1] we considered the canonical problem of assembling a peg into a hole. Simulation of this quasi-static system was used in order to select the control parameters. The integration step in the simulator was formulated as a *mixed linear complementarity problem* (MLCP). MLCPs should be thought of as *linear complementarity problems* (LCPs) coupled with additional linear equality constraints. A brief description of the linear complementarity problem and results concerning LCPs with copositive matrices are given in the following subsections. For a detailed analysis of these problems we refer the reader to the excellent manuscript [2].

1.1. Linear complementarity problems

In this section we present the definitions for the linear complementarity problem (LCP) and the mixed linear complementarity problem (MLCP).

Definition 1.1. *The problem of finding $z \in \mathbb{R}^n$ such that*

$$z \geq 0, \quad Mz + b \geq 0, \quad \text{and} \quad z^T(Mz + b) = 0, \quad (1.1)$$

where $b \in \mathbb{R}^n$ and $M \in \mathbb{R}^{n \times n}$ is called a linear complementarity problem.

In the above definition the inequality $z \geq 0$, $z \in \mathbb{R}^n$ is to be understood componentwise, i.e., $z_i \geq 0$, $i = \overline{1, n}$. The non-negativity and complementarity conditions (1.1) can be also written in the more compact form:

$$0 \leq z \perp w := Mz + b \geq 0.$$

We denote the problem (1.1) by $LCP(b, M)$. If in addition to the complementarity constraints we add some equality constraints we obtain a *mixed linear complementarity problem (MLCP)*. To be more precise, we follow the definition in [2] and consider the matrices $A \in \mathbb{R}^{n \times n}$, $B \in \mathbb{R}^{m \times m}$, $C \in \mathbb{R}^{n \times m}$ and $D \in \mathbb{R}^{m \times n}$. Let $a \in \mathbb{R}^n$ and $b \in \mathbb{R}^m$ be given.

Definition 1.2. *The mixed linear complementarity problem is the problem of finding vectors $u \in \mathbb{R}^n$ and $v \in \mathbb{R}^m$ such that*

$$\begin{aligned} a + Au + Cv &= 0 \\ b + Du + Bv &\geq 0 \\ v &\geq 0 \\ v^T(b + Du + Bv) &= 0 \end{aligned} \tag{1.2}$$

We note that if the matrix A in (1.2) is invertible we can write u in terms of v and use this form to reduce the problem to a standard LCP formulation.

1.2. LCPs with copositive matrices

The matrix of the underlying LCP used in the time-stepping schemes such as the one used in [1] is a copositive matrix.

Definition 1.3. *A matrix $M \in \mathbb{R}^{n \times n}$ is said to be copositive if*

$$x^T M x \geq 0 \text{ for all } x \in \mathbb{R}^n, x \geq 0.$$

In general a linear complementarity problem with a copositive matrix is not guaranteed to possess a solution. Solvability of such LCPs is discussed in the following Theorem.

Theorem 1.4 ([2], **Th. 3.8.6**). *Let $M \in \mathbb{R}^{n \times n}$ be a copositive matrix and let $b \in \mathbb{R}^n$ be given. If the implication*

$$[v \geq 0, Mv \geq 0, v^T M v = 0] \Rightarrow [v^T b \geq 0]$$

holds, then $LCP(b, M)$ has a solution. Lemke's algorithm with precautions taken against cycling will always find a solution of $LCP(b, M)$.

Lemke's algorithm is a pivoting method similar to the simplex method of linear programming. Cycling here refers to the possibility of using the same basis twice.

2. The quasi-static model

The continuous-time model under the rigid body assumption is given by the following *differential complementarity problem (DCP)*:

$$\dot{q}(t) = v(t), \quad (2.1)$$

$$Ev(t) - W_n(q, u, t)\lambda_n(t) - W_t(q, u, t)\lambda_t(t) = 0, \quad (2.2)$$

$$0 \leq \Psi_n(q, u, t) \perp \lambda_n(t) \geq 0, \quad (2.3)$$

$$\dot{s}_{tk}^+(t) - \dot{s}_{tk}^-(t) = (W_{tk}(q, u, t))^T v(t) + \frac{\partial \Psi_{tk}}{\partial t}(q, u, t), \quad k = 1, \dots, n_c, \quad (2.4)$$

$$0 \leq \dot{s}_{tk}^+(t) \perp \mu_k \lambda_{nk}(t) + \lambda_{tk}(t) \geq 0, \quad k = 1, \dots, n_c, \quad (2.5)$$

$$0 \leq \dot{s}_{tk}^-(t) \perp \mu_k \lambda_{nk}(t) - \lambda_{tk}(t) \geq 0, \quad k = 1, \dots, n_c. \quad (2.6)$$

Here q denotes the generalized system position and v the generalized system velocity. The control parameters are encoded in the vector u . The quasi-static assumption is reflected by the equilibrium equation (2.2), where E is a damping matrix, assumed to be symmetric positive definite. The vectors $\lambda_n(t) \in \mathbb{R}^{n_c}$ and $\lambda_t(t) \in \mathbb{R}^{n_c}$ represent all normal and tangential forces, while $W_n(q, u, t)$ and $W_t(q, u, t)$ are the normal and tangential wrench matrices. More precisely, the k -th column of $W_n(q, u, t)$ ($W_t(q, u, t)$) is the normal (tangential) wrench vector $W_{nk}(q, u, t)$ ($W_{tk}(q, u, t)$) corresponding to contact k , $k = \overline{1, n_c}$, with n_c denoting the number of active contacts. The vector $\Psi_n(q, u, t)$ contains the normal displacements for configuration q , controls u and time t . More precisely, $\Psi_n(q, u, t) = [\Psi_{n1}(q, u, t), \dots, \Psi_{nn_c}(q, u, t)]^T$, where $\Psi_{nk}(q, u, t)$ represents the normal displacement function corresponding to contact k . In a similar way, one defines the vector of tangential displacements, $\Psi_t(q, u, t) = [\Psi_{t1}(q, u, t), \dots, \Psi_{tn_c}(q, u, t)]^T$. Equation (2.3) represents the contact and non-penetration constraints; that is whenever the normal separation at contact k is strictly positive ($\Psi_{nk}(q, u, t) > 0$), the corresponding normal force is 0 ($\lambda_{nk} = 0$), while whenever contact k is established ($\Psi_{nk}(q, u, t) = 0$), the corresponding normal force is nonnegative ($\lambda_{nk} \geq 0$).

Equation (2.4) defines the positive, $\dot{s}_{tk}^+(t)$, and negative, $\dot{s}_{tk}^-(t)$, sliding velocities at contact k . The right-hand side of (2.4) represents the (overall) sliding velocity $\dot{s}_{tk}(t) := \dot{\Psi}_{tk}(q, u, t) = (W_{tk}(q, u, t))^T v(t) + \frac{\partial \Psi_{tk}}{\partial t}(q, u, t)$ at contact k . The last two equations, namely (2.5) and (2.6), represent Coulomb's friction law at contact k , with $\mu_k \in [0, 1]$ being the friction coefficients.

3. The time-stepping scheme

Let t_l denote the time at which one has a solution configuration q^l and let $t_{l+1} = t_l + h$ denote the time at which one would want an estimate of the solution. We approximate the new configuration q^{l+1} using a backward Euler formula, as follows

$$q^{l+1} = q^l + h v^{l+1},$$

where v^{l+1} is an estimate for the new velocity and will be found by solving a mixed linear complementarity problem. At each integration step the unknowns $(hv^{l+1}, h\lambda_n^{l+1}, h\lambda_f^{l+1}, h\sigma^{l+1})$ may be obtained as the solution of the following MLCP:

$$\begin{pmatrix} 0 \\ \rho_n^{l+1} \\ \rho_f^{l+1} \\ s^{l+1} \end{pmatrix} = \begin{pmatrix} E & -W_n^l & -W_f^l & 0 \\ (W_n^l)^T & 0 & 0 & 0 \\ (W_f^l)^T & 0 & 0 & E_f \\ 0 & U_f & -E_f^T & 0 \end{pmatrix} \begin{pmatrix} hv^{l+1} \\ h\lambda_n^{l+1} \\ h\lambda_f^{l+1} \\ h\sigma^{l+1} \end{pmatrix} + \begin{pmatrix} 0 \\ \Psi_n^l + h\frac{\partial\Psi_n^l}{\partial t} \\ h\frac{\partial\Psi_f^l}{\partial t} \\ 0 \end{pmatrix} \tag{3.1}$$

with $0 \leq [\rho_n^{l+1}, \rho_f^{l+1}, s^{l+1}] \perp [h\lambda_n^{l+1}, h\lambda_f^{l+1}, h\sigma^{l+1}] \geq 0$. Here $U_f \in \mathbb{R}^{n_c \times n_c}$, $E_f \in \mathbb{R}^{2n_c \times n_c}$ with U_f a diagonal matrix with elements on its diagonal equal to μ_k , $k = 1, \dots, n_c$ and E_f a block diagonal matrix, with diagonal blocks given by the vector e (e is a two-dimensional vector of all ones). That is,

$$U_f = \begin{pmatrix} \mu_1 & \dots & 0 \\ \vdots & & \vdots \\ 0 & \dots & \mu_{n_c} \end{pmatrix}, \quad E_f = \begin{pmatrix} 1 & \dots & 0 \\ 1 & \dots & 0 \\ \vdots & & \vdots \\ 0 & \dots & 1 \\ 0 & \dots & 1 \end{pmatrix}.$$

The superscript l used in the MLCP (3.1) indicates that all the corresponding quantities are calculated with $q := q^l$ and $t := t_l$. For each contact k we define the 3×2 matrix $W_{fk}(q, u, t)$ by joining the column vectors $W_{tk}(q, u, t)$ and $-W_{tk}(q, u, t)$. That is,

$$W_{fk}(q, u, t) = [W_{tk}(q, u, t) \quad -W_{tk}(q, u, t)].$$

If we put all the active contacts together we obtain the "frictional" wrench matrix $W_f(q, u, t)$ appearing in formulation (3.1). In a similar way, we get the vector $\Psi_f(q, u, t)$.

Solvability and the Friction Cone. For an active contact k , we define the friction cone corresponding to that contact by

$$FC_k(q, u, t) = \{z = W_{nk}\lambda_{nk} + W_{fk}\lambda_{fk} \mid \lambda_{nk} \geq 0, \lambda_{fk} \geq 0, e^T \lambda_{fk} \leq \mu_k \lambda_{nk}\}, \tag{3.2}$$

where $W_{nk} := W_{n,k}(q, u, t)$, $W_{fk} := W_{f,k}(q, u, t)$ and $e = [1, 1]^T$. The total friction cone, $FC(q, u, t)$, which accounts for all active contacts is defined by

$$FC(q, u, t) = \sum_{k=1}^{n_c} FC_k(q, u, t).$$

Using the fact that the matrix E in the MLCP (3.1) is positive definite, we can eliminate the variables hv^{l+1} and reduce the MLCP to a standard LCP with a copositive matrix. It can be shown that the resulting LCP, is solvable whenever the total friction cone $FC(q^l, u, t_l)$ is pointed. We recall that a cone is pointed if it doesn't contain any proper subspace. The lack of pointedness for the friction cone results in jammed configurations (see [3]) and therefore this regularity assumption is very realistic and can be successfully used in devising randomized plans (see [1]).

4. Conclusions

We have discussed an LCP-based time-stepping scheme that can be used to simulate rigid body systems in a quasi-static setting. The scheme was introduced and successfully used for a particular case in [1]. Solvability of the integration step is guaranteed by the pointedness of the friction cone, an assumption that is common in dynamic settings as well (see [3] and [4] for example).

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