

## PFAFFIAN TRANSFORMATIONS

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**Abstract.** Geometrical and structural properties are proved for manifolds possessing a particular locally conformal almost cosymplectic structure.

### 1. Introduction

Let  $M(g, \Omega, \phi, \eta, \xi)$  be an  $2m + 1$ -dimensional Riemannian manifold with metric tensor  $g$  and associated Levi-Civita connection  $\nabla$ . The quadruple  $(\Omega, \phi, \xi, \eta)$  consists of a structure 2-form  $\Omega$  of rank  $2m$ , an endomorphism  $\phi$  of the tangent bundle, the Reeb vector field  $\xi$ , and its corresponding Reeb covector field  $\eta$ , respectively.

We assume that the 2-form  $\Omega$  satisfies the relation

$$d\Omega = \lambda \eta \wedge \Omega, \tag{1}$$

where  $\lambda$  is constant, and that the 1-form  $\eta$  is given by

$$\eta = \lambda df, \tag{2}$$

for some scalar function  $f$  on  $M$ . We may therefore notice that a locally conformal almost cosymplectic structure [7] [10] is defined on the manifold  $M$ .

In addition, we assume that the field  $\phi$  of endomorphisms of the tangent spaces defines a quasi-Sasakian structure, thus realizing in particular the identity

$$\phi^2 = -\text{Id} + \eta \otimes \xi.$$

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Moreover, we will assume the presence on  $M$  of a structure vector field  $X$  satisfying the property

$$\nabla X = fdp + \lambda \nabla \xi. \quad (3)$$

In the present paper various properties involving the above mentioned objects are studied. In particular, for the Lie differential of  $\Omega$  and  $\eta$  with respect to  $X$ , one has

$$\begin{aligned} \mathcal{L}_X \eta &= 0, \\ \mathcal{L}_X \Omega &= 0, \end{aligned}$$

which shows that  $\eta$  and  $\Omega$  define Pfaffian transformations [3].

## 2. Preliminaries

Let  $(M, g)$  be an  $n$ -dimensional Riemannian manifold and let  $\nabla$  be the covariant differential operator defined by the metric tensor. We assume in the sequel that  $M$  is oriented and that the connection  $\nabla$  is symmetric.

Let  $\Gamma TM = \Xi(M)$  be the set of sections of the tangent bundle  $TM$ , and

$$\flat : TM \xrightarrow{\flat} T^*M \quad \text{and} \quad \sharp : T^*M \xleftarrow{\sharp} TM$$

the classical isomorphisms defined by the metric tensor  $g$  (i.e.  $\flat$  is the index lowering operator, and  $\sharp$  is the index raising operator).

Following [12], we denote by

$$A^q(M, TM) = \Gamma \text{Hom}(\Lambda^q TM, TM),$$

the set of vector valued  $q$ -forms ( $q < \dim M$ ), and we write for the covariant derivative operator with respect to  $\nabla$

$$d^\nabla : A^q(M, TM) \rightarrow A^{q+1}(M, TM). \quad (4)$$

It should be noticed that in general  $d^{\nabla^2} = d^\nabla \circ d^\nabla \neq 0$ , unlike  $d^2 = d \circ d = 0$ .

Furthermore, we denote by  $dp \in A^1(M, TM)$  the canonical vector valued 1-form of  $M$ , which is also called the soldering form of  $M$  [3]; since  $\nabla$  is assumed to be symmetric, we recall that the identity  $d^\nabla(dp) = 0$  is valid.

The operator

$$d^\omega = d + e(\omega),$$

acting on  $\Lambda M$  is called the cohomology operator [5]. Here,  $e(\omega)$  means the exterior product by the closed 1-form  $\omega$ , i.e.

$$d^\omega u = du + \omega \wedge u,$$

with  $u \in \Lambda M$ . A form  $u \in \Lambda M$  such that

$$d^\omega u = 0,$$

is said to be  $d^\omega$ -closed, and  $\omega$  is called the cohomology form.

A vector field  $X \in \Xi(M)$  which satisfies

$$d^\nabla(\nabla X) = \nabla^2 X = \pi \wedge dp \in A^2(M, TM), \quad \pi \in \Lambda^1 M, \quad (5)$$

and where  $\pi$  is conformal to  $X^\flat$ , is defined to be an exterior concurrent vector field [14]. In this case, if  $\mathcal{R}$  denotes the Ricci tensor field of  $\nabla$ , one has

$$\mathcal{R}(X, Z) = -2m\lambda^3(\kappa + \eta) \wedge dp, \quad Z \in \Xi(M)$$

### 3. Geometrical properties

In terms of a local field of adapted vectorial frames  $\mathcal{O} = \text{vect}\{e_A | A = 0, \dots, 2m\}$  and its associated coframe  $\mathcal{O}^* = \text{covect}\{\omega^A | A = 0, \dots, 2m\}$ , the soldering form  $dp$  can be expressed as

$$dp = \sum_{A=0}^{2m} \omega^A \otimes e_A;$$

and we recall that E. Cartan's structure equations can be written as

$$\nabla e_A = \sum_{B=0}^{2m} \theta_A^B \otimes e_B, \quad (6)$$

$$d\omega^A = - \sum_{B=0}^{2m} \theta_B^A \wedge \omega^B, \quad (7)$$

$$d\theta_B^A = - \sum_{C=0}^{2m} \theta_B^C \wedge \theta_C^A + \Theta_B^A. \quad (8)$$

In the above equations  $\theta$  (respectively  $\Theta$ ) are the local connection forms in the tangent bundle  $TM$  (respectively the curvature 2-forms on  $M$ ).

In terms of the frame fields  $\mathcal{O}$  and  $\mathcal{O}^*$  with  $e_0 = \xi$  and  $\omega^0 = \eta$ , the structure vector field  $X$  and the 2-form  $\Omega$  can be expressed as

$$X = \sum_{a=1}^{2m} X^a e_a, \quad (9)$$

$$\Omega = \sum_{i=1}^m \omega^i \wedge \omega^{i^*}, \quad i^* = i + m. \quad (10)$$

Taking the Lie differential of  $\Omega$  and  $\eta$  with respect to  $X$ , one calculates

$$\mathcal{L}_X \eta = 0, \quad (11)$$

$$\mathcal{L}_X \Omega = 0. \quad (12)$$

According to [6] the above equations (11) and (12) prove that that  $\eta$  and  $\Omega$  define a Pfaffian transformation [3].

Next, by (2) one gets that

$$\theta_0^a = \lambda \omega^a. \quad (13)$$

Since we also assume that

$$\nabla X = f dp + \lambda \nabla \xi, \quad (14)$$

we further also derive that

$$\nabla \xi = \lambda(dp - \eta \otimes \xi). \quad (15)$$

Since the  $q$ -th covariant differential  $\nabla^q Z$  of a vector field  $Z \in \Xi(M)$  is defined inductively, i.e.

$$\nabla^q Z = d^\nabla(\nabla^{q-1} Z),$$

this yields

$$\nabla^2 \xi = \lambda^2 \eta \otimes dp, \quad (16)$$

$$\nabla^3 \xi = 0. \quad (17)$$

Hence, one may say that the 3-covariant Reeb vector field  $\xi$  is vanishing.

Next, by (13), one derives that

$$\nabla^2 X = \lambda^3 (df + \eta) \wedge dp = \frac{1 + \lambda}{\lambda} \eta \wedge dp, \quad (18)$$

and consecutively one gets that

$$\nabla^3 X = 0. \quad (19)$$

This shows that both vector fields  $\xi$  and  $X$  together define a 3-vanishing structure.

Moreover, by reference to [13], it follows from (18) that one may write that

$$\nabla^2 X = -\frac{1}{2m} \text{Ric}(X) - X^\flat \wedge dp, \quad (20)$$

where Ric is the Ricci tensor.

Reminding that by the definition of the operator  $\phi$

$$\begin{aligned} \phi e_i &= e_{i^*} & i &\in \{1, \dots, m\}, \\ \phi e_{i^*} &= -e_i & i^* &= i + m, \end{aligned}$$

one can check that indeed  $\phi^2 = -\text{Id}$ . Acting with  $\phi$  on the vector field  $X$ , one obtains in a first step that

$$\phi X = \sum_{i=1}^m X^i e_{i^*} - X^{i^*} e_i \quad i^* = i + m. \quad (21)$$

Calculating the Lie derivative of  $\phi$  w.r.t.  $\xi$ , one gets

$$(\mathcal{L}_\xi \phi)X = [\xi, \phi X] - \phi[\xi, X]. \quad (22)$$

Since clearly

$$[\xi, \phi X] = 0, \quad (23)$$

there follows that

$$(\mathcal{L}_\xi \phi)X = 0. \quad (24)$$

Hence, the Jacobi bracket corresponding to the Reeb vector field  $\xi$  vanishes.

By reference to the definition of the divergence

$$\operatorname{div} Z = \sum_{A=0}^{2m} \omega^A (\nabla_{e_A} Z)$$

one obtains in the case under consideration that

$$\operatorname{div} X = 2m(\lambda + f^2), \quad (25)$$

and

$$\operatorname{div} \phi X = 0. \quad (26)$$

Calculating the differential of the dual form  $X^\flat$  of  $X$ , one gets

$$dX^\flat = \sum_{a=1}^{2m} \left( dX^a + \sum_{b=1}^{2m} X^b \theta_b^a \right) \wedge \omega^a. \quad (27)$$

Since

$$dX^a + \sum_{b=1}^{2m} X^b \theta_b^a = \lambda \omega^a, \quad (28)$$

one has that

$$dX^\flat = 0, \quad (29)$$

which means that the Pfaffian  $X^\flat$  is closed. This implies that  $X^\flat$  is an eigenfunction of the Laplacian  $\Delta$ , and one can write that

$$\Delta X^\flat = f \|X\|^2 X^\flat.$$

If we set

$$2l = \|X\|^2, \quad (30)$$

one also derives by (28) that

$$dl = \lambda X^b. \quad (31)$$

From (31) it follows that  $dX^b = 0$  which is indeed in accordance with (29).

Returning to the operator  $\phi$ , one calculates that

$$\nabla(\phi X) = \lambda \phi dp - \sum_{i=1}^m \left( \sum_{a=1}^{2m} (X^a \theta_a^i) \otimes e_{i^*} + \sum_{a=1}^{2m} (X^a \theta_a^{i^*}) \otimes e_i \right). \quad (32)$$

Hence there follows that

$$[\xi, X] = \rho \xi - \phi C, \quad (33)$$

$$[\xi, \phi X] = ((C^0)^2 + C^0(1 - \lambda))\xi, \quad (34)$$

$$[X, \phi X] = \nabla_\xi \phi C = C^0 \xi - C \quad (35)$$

which shows that the triple  $\{X, \xi, \phi X\}$  defines a 3-distribution on  $M$ .

It is also interesting to draw the attention on the fact that  $X$  possesses the following property. From (14) and (15) one derives that

$$\nabla_X X = f X, \quad (36)$$

which means that  $X$  is an affine geodesic vector field.

Finally, if we denote by  $\Sigma$  the exterior differential system which defines  $X$ , it follows by Cartan's test [1] that the characteristic numbers are

$$r = 3, \quad s_0 = 1, \quad s_1 = 2.$$

Since  $r = s_0 + s_1$ , it follows that  $\Sigma$  is in involution and the existence of  $X$  depends on an arbitrary function of 1 argument.

Summarizing, we can organize our results into the following

**Theorem 3.1.** *Let  $M$  be a  $2m + 1$ -dimensional Riemannian manifold and let  $\nabla$  be the Levi-Civita connection and  $\xi$  be the Reeb vector field and  $\eta$  the Reeb covector field on  $M$ . One has the following properties:*

- (i):  $\xi$  and  $X$  define a 3-vanishing structure;
- (ii): the Jacobi bracket corresponding to  $\xi$  vanishes;
- (iii): the harmonic operator acting on  $X^\flat$  gives

$$\Delta X^\flat = f \|X\|^2 X^\flat,$$

which proves that  $X^\flat$  is an eigenfunction of  $\Delta$ , having  $f \|X\|^2$  as eigenvalue;

- (iv): the 2-form  $\Omega$  and the Reeb covector  $\eta$  define a Pfaffian transformation, i.e.

$$\begin{aligned} \mathcal{L}_X \Omega &= 0, \\ \mathcal{L}_X \eta &= 0; \end{aligned}$$

- (v): the Ricci tensor is determined by  $\nabla^2 X$ ;
- (vi): one has

$$\nabla_X X = fX, \quad f = \text{scalar},$$

which shows that  $X$  is an affine geodesic;

- (vii): the triple  $\{X, \xi, \phi X\}$  is a 3-distribution on  $M$  and is in involution in the sense of Cartan.

#### 4. The structure 2-form $\Omega$

In the present section, we derive some properties of the structure 2-form  $\Omega$ . First, we recall that one has

$$d\Omega = \lambda \eta \wedge \Omega, \quad \lambda = \text{constant}. \quad (37)$$

By Lie differentiation with respect to  $X$ , one gets

$$\mathcal{L}_X \Omega = 0. \quad (38)$$



Further, since  $i_\xi \Omega = 0$ , one calculates that

$$\begin{aligned}\mathcal{L}_\xi \Omega &= \lambda \Omega, \\ d(\mathcal{L}_\xi \Omega) &= \lambda^2 \eta \wedge \Omega.\end{aligned}$$

Moreover, by the Lie bracket  $[ \cdot, \cdot ]$  one also has that

$$i_{[X, \xi]} \Omega = 0. \quad (39)$$

Next, we consider the vector field  $\phi X$ . By (32), one calculates that

$$\mathcal{L}_{\phi X} \Omega = -2\lambda \eta \wedge X^\flat, \quad \lambda = \text{constant}. \quad (40)$$

Since  $X^\flat$  is closed, this yields

$$d(\mathcal{L}_{\phi X} \Omega) = 0. \quad (41)$$

This shows that  $\phi X$  defines a relative conformal transformation [15] [8] of  $\Omega$ . In addition, one also derives that

$$\mathcal{L}_{[X, \xi]} \Omega = \mathcal{L}_X \mathcal{L}_\xi \Omega - \mathcal{L}_\xi \mathcal{L}_X \Omega = \mathcal{L}_X \mathcal{L}_\xi \Omega$$

and

$$\mathcal{L}_{fX} \Omega = f \mathcal{L}_X \Omega + df \wedge i_X \Omega = df \wedge i_X \Omega$$

**Theorem 4.1.** *The structure 2-form  $\Omega$  satisfies the following relations*

<p>(i):</p> $d\Omega = \lambda \eta \wedge \Omega$	<p>(iv):</p> $d(\mathcal{L}_{\phi X} \Omega) = 0$
<p>(ii):</p> $\mathcal{L}_\xi \Omega = \lambda \Omega$ $d(\mathcal{L}_\xi \Omega) = \lambda^2 \eta \wedge \Omega$	<p>(v):</p> $\mathcal{L}_{[X, \xi]} \Omega = \mathcal{L}_X \mathcal{L}_\xi \Omega$
<p>(iii):</p> $i_{[X, \xi]} \Omega = 0$	<p>(vi):</p> $\mathcal{L}_{fX} \Omega = df \wedge i_X \Omega$

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