# A NOTE ON THREE-STEP ITERATIVE METHODS FOR NONLINEAR EQUATIONS

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Dedicated to Professor Petru Blaga at his 60<sup>th</sup> anniversary

**Abstract**. In this short note we give certain comments and improvements of some three-step iterative methods recently considered by N.A. Mir and T. Zaman (Appl. Math. Comput. (2007) doi: 10.1016/j.amc.2007.03.071).

# 1. Introduction

Very recently, N.A. Mir and T. Zaman [1] have considered three-step quadrature based iterative methods for finding a single zero  $x = \alpha$  of a nonlinear equation

$$f(x) = 0. (1.1)$$

All variants of their methods include the formula

$$x_{n+1} = x_n - \frac{f(x_n)f'(x_n)}{f'(x_n)^2 - \lambda f(x_n)f''(x_n)},$$
(1.2)

obtained from the rectangular quadrature formula. It is clear that (1.2) reduces to Newton and Halley method for  $\lambda = 0$  and  $\lambda = 1/2$ , respectively.

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As a variant with maximal order of convergence they have proposed the following three-step method

$$y_{n} = x_{n} - \frac{f(x_{n})}{f'(x_{n})},$$

$$z_{n} = y_{n} - \frac{f(y_{n})f'(y_{n})}{f'(y_{n})^{2} - \lambda f(y_{n})f''(y_{n})},$$

$$x_{n+1} = z_{n} - \frac{(y_{n} - z_{n})f(z_{n})}{f(y_{n}) - 2f(z_{n})},$$

$$(1.3)$$

proving that for a sufficiently smooth function f and a starting point  $x_0$  sufficiently close to the single zero  $x = \alpha$ , this method has eighth order convergence for  $\lambda = 1/2$ , i.e.,

$$e_{n+1} = (-c_3c_2^5 + c_2^7)e_n^8 + O(e_n^9),$$

where  $e_n = x_n - \alpha$  and

$$c_k = \frac{1}{k!} \frac{f^{(k)}(\alpha)}{f'(\alpha)}, \quad k = 2, 3, \dots$$
 (1.4)

As we can see this three-step method need six function evaluations per iteration:  $f(x_n)$ ,  $f(y_n)$ ,  $f(z_n)$ ,  $f'(x_n)$ ,  $f'(y_n)$ , and  $f''(y_n)$ . Without new function evaluations, in this note we show that the formula

$$x_{n+1} = S(y_n, z_n) = z_n - \frac{f(z_n)}{f'(y_n) + (z_n - y_n)f''(y_n)}$$
(1.5)

is a much better choice than the third formula in (1.3). In that case the corresponding three-step method has tenth order convergence. Moreover, the formula (1.5) is numerically stable in comparing with the previous one.

The paper is organized as follows. In Section 2 we give certain auxiliary formulae, which can be used also in other investigations in convergence analysis. The main results and a numerical example are given in Section 3.

# 2. Some auxiliary formulae

We suppose that the equation (1.1) has a single zero  $x = \alpha$  in certain neighborhood  $U_{\varepsilon}(\alpha) := (\alpha - \varepsilon, \alpha + \varepsilon), \ \varepsilon > 0$ , and that the function f is sufficiently differentiable in  $U_{\varepsilon}(\alpha)$ . Evidently,  $f'(\alpha) \neq 0$ .

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Let  $x_n \in U_{\varepsilon}(\alpha)$  and

$$e_n := x_n - \alpha, \quad \widetilde{e}_n := y_n - \alpha, \quad \widehat{e}_n := z_n - \alpha.$$

Using (1.4) it is easy to get the following formula

$$\frac{f(x_n)}{f'(x_n)} = e_n - c_2 e_n^2 + 2(c_2^2 - c_3) e_n^3 - (4c_2^3 - 7c_3c_2 + 3c_4) e_n^4 
+ (8c_2^4 - 20c_3c_2^2 + 10c_4c_2 + 6c_3^2 - 4c_5) e_n^5 
- [16c_2^5 - 52c_3c_2^3 + 28c_4c_2^2 + (33c_3^2 - 13c_5)c_2 - 17c_3c_4 + 5c_6] e_n^6 
+ O(e_n^7).$$
(2.1)

This formula is an inverse of the well-known Schröder formula (cf. [2, pp. 352–354]).

Therefore, in the case of the Newton method

$$\Phi_N(x_n) = x_n - \frac{f(x_n)}{f'(x_n)} \tag{2.2}$$

we have

$$\widetilde{e}_{n} = \Phi_{N}(x_{n}) - \alpha$$

$$= c_{2}e_{n}^{2} - 2(c_{2}^{2} - c_{3})e_{n}^{3} + (4c_{2}^{3} - 7c_{3}c_{2} + 3c_{4})e_{n}^{4}$$

$$-(8c_{2}^{4} - 20c_{3}c_{2}^{2} + 10c_{4}c_{2} + 6c_{3}^{2} - 4c_{5})e_{n}^{5}$$

$$+[16c_{2}^{5} - 52c_{3}c_{2}^{3} + 28c_{4}c_{2}^{2} + (33c_{3}^{2} - 13c_{5})c_{2} - 17c_{3}c_{4} + 5c_{6}]e_{n}^{6}$$

$$+O(e_{n}^{7}).$$
(2.3)

Also, we need the corresponding expression for

$$\widetilde{C}_{2}(y_{n}) := \frac{f''(y_{n})}{2f'(y_{n})} = \frac{1}{2} \frac{f''(\alpha) + \frac{f'''(\alpha)}{1!} \widetilde{e}_{n} + \frac{f^{iv}(\alpha)}{2!} \widetilde{e}_{n}^{2} + \cdots}{f'(\alpha) + \frac{f'''(\alpha)}{1!} \widetilde{e}_{n} + \frac{f'''(\alpha)}{2!} \widetilde{e}_{n}^{2} + \cdots},$$

i.e.,

$$\widetilde{C}_2(y_n) = \frac{1}{2} \cdot \frac{1 \cdot 2 c_2 + 2 \cdot 3 c_3 \,\widetilde{e}_n + 3 \cdot 4 c_4 \,\widetilde{e}_n^2 + \cdots}{1 + 2 c_2 \,\widetilde{e}_n + 3 c_3 \,\widetilde{e}_n^2 + \cdots},$$

where  $c_k$  are defined by (1.4). It gives

$$\widetilde{C}_2(y_n) = A_0 + A_1 \widetilde{e}_n + A_2 \widetilde{e}_n^2 + A_3 \widetilde{e}_n^3 + \cdots,$$
 (2.4)

where

$$A_0 = c_2, \quad A_1 = 3c_3 - 2c_2^2, \quad A_2 = 4c_2^3 - 9c_3c_2 + 6c_4,$$

$$A_3 = -8c_2^4 + 24c_3c_2^2 - 16c_4c_2 - 9c_3^2 + 10c_5,$$

$$A_4 = 16c_2^5 - 60c_3c_3^3 + 40c_4c_2^2 + 5(9c_3^2 - 5c_5)c_2 + 15(c_6 - 2c_3c_4),$$

$$A_5 = -32c_2^6 + 144c_3c_2^4 - 96c_4c_2^3 + (60c_5 - 162c_3^2)c_2^2 + 36(4c_3c_4 - c_6)c_2$$

$$+3(9c_3^3 - 15c_5c_3 - 8c_4^2 + 7c_7),$$

$$A_6 = 64c_2^7 - 336c_3c_2^5 + 224c_4c_2^4 + 28(18c_3^2 - 5c_5)c_2^3 - 84(6c_3c_4 - c_6)c_2^2$$

$$-7(27c_3^3 - 30c_5c_3 - 16c_4^2 + 7c_7)c_2 + 7(18c_4c_3^2 - 9c_6c_3 - 10c_4c_5 + 4c_8),$$

etc.

Now, for the Halley method

$$\Phi_H(y_n) = y_n - \frac{f(y_n)/f'(y_n)}{1 - \tilde{C}_2(y_n)(f(y_n)/f'(y_n))}$$
(2.5)

we have

$$\hat{e}_n = \Phi_H(y_n) - \alpha = \tilde{e}_n - g_n \left( 1 + \tilde{C}_2(y_n)g_n + [\tilde{C}_2(y_n)g_n]^2 + \cdots \right), \quad g_n = \frac{f(y_n)}{f'(y_n)}.$$

Using (2.1), in this case, we get

$$\hat{e}_n = (c_2^2 - c_3)\tilde{e}_n^3 - 3(c_2^3 - 2c_3c_2 + c_4)\tilde{e}_n^4 + 6(c_2^4 - 3c_3c_2^2 + 2c_4c_2 + c_3^2 - c_5)\tilde{e}_n^5$$

$$-[9c_2^5 - 37c_3c_2^3 + 29c_4c_2^2 + 4(7c_3^2 - 5c_5)c_2 - 19c_3c_4 + 10c_6]\tilde{e}_n^6 + O(\tilde{e}_n^7).$$
(2.6)

In our analysis we also need an expansion of  $f(z_n)/f'(y_n)$  in terms of  $\tilde{e}_n$  (=  $y_n - \alpha$ ), where  $z_n - \alpha = \hat{e}_n$  is given by (2.6). Thus, we have

$$v_n = \frac{f(z_n)}{f'(y_n)} = \frac{\hat{e}_n + c_2 \,\hat{e}_n^2 + c_3 \,\hat{e}_n^3 + \cdots}{1 + 2c_2 \,\tilde{e}_n + 3c_3 \,\tilde{e}_n^2 + \cdots},$$

i.e.,

$$v_{n} = (c_{2}^{2} - c_{3})\tilde{e}_{n}^{3} - (5c_{2}^{3} - 8c_{3}c_{2} + 3c_{4})\tilde{e}_{n}^{4}$$

$$+ (16c_{2}^{4} - 37c_{3}c_{2}^{2} + 18c_{4}c_{2} + 9c_{3}^{2} - 6c_{5})\tilde{e}_{n}^{5}$$

$$- (40c_{2}^{5} - 124c_{3}c_{2}^{3} + 69c_{4}c_{2}^{2} - (32c_{5} - 69c_{3}^{2})c_{2} - 32c_{3}c_{4} + 10c_{6})\tilde{e}_{n}^{6}$$

$$+ O(\tilde{e}_{n}^{7}). \tag{2.7}$$

In the case when  $y_n = \Phi_N(x_n)$  and

$$z_n = y_n - \frac{(x_n - y_n)f(y_n)}{f(x_n) - 2f(y_n)},$$
(2.8)

we are interested in

$$u_n = \frac{f(x_n)}{f(y_n)}, \quad t_n = \frac{f(x_n)}{f'(z_n)}, \quad s_n = \frac{f'(x_n)}{f'(z_n)},$$
 (2.9)

i.e.,

$$u_n = \frac{e_n + c_2 e_n^2 + c_3 e_n^3 + \cdots}{\tilde{e}_n + c_2 \tilde{e}_n^2 + c_3 \tilde{e}_n^3 + \cdots},$$
(2.10)

$$t_n = \frac{e_n + c_2 e_n^2 + c_3 e_n^3 + \cdots}{1 + 2c_2 \bar{e}_n + 3c_3 \bar{e}_n^2 + \cdots}, \quad \text{and} \quad s_n = \frac{1 + 2c_2 e_n + 3c_3 e_n^2 + \cdots}{1 + 2c_2 \bar{e}_n + 3c_3 \bar{e}_n^2 + \cdots},$$

where  $e_n = x_n - \alpha$ ,  $\tilde{e}_n = y_n - \alpha$ , and  $\bar{e}_n = z_n - \alpha$ .

Here,  $\bar{e}_n = \tilde{e}_n - (e_n - \tilde{e}_n)/(u_n - 2)$ . According to (2.3) and (2.10) we get

$$\bar{e}_n = c_2(c_2^2 - c_3)e_n^4 - 2(2c_2^4 - 4c_3c_2^2 + c_4c_2 + c_3^2)e_n^5$$

$$+ [10c_2^5 - 30c_3c_2^3 + 12c_4c_2^2 + 3c_2(6c_3^2 - c_5) - 7c_3c_4]e_n^6$$

$$+ O(e_n^7).$$

$$(2.11)$$

For  $t_n$  and  $s_n$  we obtain

$$t_n = e_n + c_2 e_n^2 + c_3 e_n^3 + c_4 e_n^4 - (2c_2^4 - 2c_3 c_2^2 - c_5) e_n^5$$

$$+ (6c_2^5 - 14c_3 c_2^3 + 4c_4 c_2^2 + 4c_3^2 c_2 + c_6) e_n^6 + O(e_n^7)$$
(2.12)

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and

$$s_{n} = 1 + 2c_{2}e_{n} + 3c_{3}e_{n}^{2} + 4c_{4}e_{n}^{3} - (2c_{2}^{4} - 2c_{3}c_{2}^{2} - 5c_{5})e_{n}^{4}$$

$$+ (4c_{2}^{5} - 12c_{3}c_{2}^{3} + 4c_{4}c_{2}^{2} + 4c_{3}^{2}c_{2} + 6c_{6})e_{n}^{5}$$

$$- [4c_{2}^{6} - 22c_{3}c_{2}^{4} + 16c_{4}c_{2}^{3} - (6c_{5} - 22c_{3}^{2})c_{2}^{2} - 14c_{3}c_{4}c_{2} - 7c_{7}]e_{n}^{6}$$

$$+ O(e_{n}^{7}), \qquad (2.13)$$

respectively.

## 3. Main results

We consider now the third-step iterative formula given by (2.2), (2.5), and (1.5), i.e.,

$$y_n = \Phi_N(x_n), \quad z_n = \Phi_H(y_n), \quad e_{n+1} = S(y_n, z_n), \quad n = 0, 1, \dots,$$
 (3.1)

for finding a simple zero  $x = \alpha$  of the equation (1.1).

**Theorem 3.1.** For a sufficiently differentiable function f in  $U_{\varepsilon}(\alpha)$  and  $x_0$  sufficiently close to  $\alpha$ , the third-step method (3.1) has tenth order of convergence, i.e.,

$$e_{n+1} = 3c_2^5 c_3(c_3 - c_2^2)e_n^{10} + 30c_2^4(c_2^2 - c_3)^2 c_3 e_n^{11} + O(e_n^{12})$$
(3.2)

where  $e_n = x_n - \alpha$  and  $c_k$  are given in (1.4).

*Proof.* According to (3.1) (and (1.5)) we have

$$e_{n+1} = x_{n+1} - \alpha = S(y_n, z_n) - \alpha = \hat{e}_n - \frac{f(z_n)/f'(y_n)}{1 + 2(\hat{e}_n - \tilde{e}_n)\tilde{C}_2(y_n)},$$

where  $\tilde{e}_n = y_n - \alpha$ ,  $\hat{e}_n = z_n - \alpha$ , and  $\tilde{C}_2(y_n) = f''(y_n)/(2f'(y_n))$ . Replacing  $\tilde{C}_2(y_n)$  and  $v_n = f(z_n)/f'(y_n)$  by the corresponding expressions (2.4) and (2.7), we obtain

$$e_{n+1} = 3c_3(c_3 - c_2^2)\widetilde{e}_n^5 + (c_2^5 + 7c_3c_2^3 - 8c_4c_2^2 - 17c_3^2c_2 + 17c_3c_4)\widetilde{e}_n^6 + O(\widetilde{e}_n^7).$$

Finally, using (2.3) we get (3.2).  $\square$ 

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In [1] the authors also considered the following three-step method

$$y_{n} = x_{n} - \frac{f(x_{n})}{f'(x_{n})},$$

$$z_{n} = y_{n} - \frac{(x_{n} - y_{n})f(y_{n})}{f(x_{n}) - 2f(y_{n})},$$

$$x_{n+1} = z_{n} - \frac{f(z_{n})f'(z_{n})}{f'(z_{n})^{2} - \lambda f(z_{n}) \left\{ \frac{2(f(z_{n}) - f(x_{n}))}{(z_{n} - x_{n})^{2}} - \frac{2f'(x_{n})}{z_{n} - x_{n}} \right\}},$$
(3.3)

with five function evaluations per iteration:  $f(x_n)$ ,  $f(y_n)$ ,  $f(z_n)$ ,  $f'(x_n)$ , and  $f'(z_n)$ . Their Theorem 3 states that this method has seventh order of convergence for any value of  $\lambda$ . However, the order of convergence is bigger than seven. Namely, we have the following result:

**Theorem 3.2.** For a sufficiently differentiable function f in  $U_{\varepsilon}(\alpha)$  and  $x_0$  sufficiently close to  $\alpha$ , the third-step method (3.3) has eighth order of convergence for any  $\lambda \neq 1/2$ , except for  $\lambda = 1/2$  when the convergence is of the order nine. Then,

$$e_{n+1} = -2c_3c_2^2(c_2^2 - c_3)^2 e_n^9$$

$$+c_2(c_2^2 - c_3)(16c_3c_2^4 - 3c_4c_2^3 - 32c_3^2c_2^2 + 11c_2c_3c_4 + 8c_3^3)e_n^{10}$$

$$+O(e_n^{11}),$$

$$(3.4)$$

where  $e_n = x_n - \alpha$  and  $c_k$  are given in (1.4).

*Proof.* Using the expansion (2.1) for the Newton correction  $f(x_n)/f'(x_n) =$ :  $h(e_n)$ , we have  $f(z_n)/f'(z_n) = h(\bar{e}_n)$ , where  $\bar{e}_n$  is given by (2.11). According to (2.9), for the third formula in (3.3) we get

$$e_{n+1} = \bar{e}_n - \frac{h(\bar{e}_n)}{1 - 2\lambda \left\{ \frac{h(\bar{e}_n) - t_n}{(\bar{e}_n - e_n)^2} - \frac{s_n}{\bar{e}_n - e_n} \right\}},$$
(3.5)

where the expansions for  $t_n$  and  $s_n$  are given by (2.12) and (2.13), respectively. This gives

$$e_{n+1} = (1 - 2\lambda)c_2^3(c_2^2 - c_3)^2 e_n^8 + 4c_2^2(c_2^2 - c_3) [2(2\lambda - 1)c_2^4 + (4 - 9\lambda)c_3c_2^2 + (2\lambda - 1)c_4c_2 + (3\lambda - 1)c_3^2]e_n^9 + \cdots$$

For  $\lambda = 1/2$  it reduces to (3.4).

Thus, the computational efficiency of the method (3.3), for  $\lambda = 1/2$ , is EFF=  $9^{1/5} \approx 1.55185$ . With the same function evaluations we can get a slightly simpler method of order nine with the same efficiency.

**Theorem 3.3.** For a sufficiently differentiable function f in  $U_{\varepsilon}(\alpha)$  and  $x_0$  sufficiently close to  $\alpha$ , the third-step method

$$y_{n} = x_{n} - \frac{f(x_{n})}{f'(x_{n})},$$

$$z_{n} = y_{n} - \frac{(x_{n} - y_{n})f(y_{n})}{f(x_{n}) - 2f(y_{n})},$$

$$x_{n+1} = z_{n} - \frac{f(z_{n})f'(z_{n})}{f'(z_{n})^{2} - \frac{1}{2}f(z_{n})\frac{f'(z_{n}) - f'(x_{n})}{z_{n} - x_{n}}},$$

$$(3.6)$$

has ninth order of convergence, i.e.,

$$e_{n+1} = -\frac{3}{2}c_3c_2^2(c_2^2 - c_3)^2 e_n^9$$

$$+2c_2(c_2^2 - c_3)(6c_2^4c_3 - 12c_2^2c_3^2 + 3c_3^3 - c_2^3c_4 + 4c_2c_3c_4)e_n^{10}$$

$$+O(e_n^{11}), \tag{3.7}$$

where  $e_n = x_n - \alpha$  and  $c_k$  are given in (1.4).

*Proof.* Similarly as in the proof of the previous theorem, we have now

$$e_{n+1} = \bar{e}_n - \frac{h(\bar{e}_n)}{1 - \frac{1}{2}h(\bar{e}_n)\frac{1 - s_n}{\bar{e}_n - e_n}}$$

instead of (3.5). This gives (3.7).

The number of function evaluations in (3.6) can be reduced to four if we take an approximation of  $f'(z_n)$  in the form

$$f'(z_n) \approx \tilde{f}'(z_n) = p_n f(x_n) + q_n f(y_n) + r_n f(z_n) + w_n f'(x_n),$$

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obtained by the Hermite interpolation (cf. [3, pp. 51–58]), where

$$p_n = \frac{(y_n - z_n)(z_n + 2y_n - 3x_n)}{(x_n - y_n)^2(x_n - z_n)}, \quad q_n = \frac{(x_n - z_n)^2}{(x_n - y_n)^2(y_n - z_n)},$$

$$r_n = \frac{3z_n - 2y_n - x_n}{(x_n - z_n)(y_n - z_n)}, \qquad w_n = \frac{y_n - z_n}{x_n - y_n}.$$

For a such modified three-step method, in notation  $(3.6^{M})$ , the following result holds:

**Theorem 3.4.** For a sufficiently differentiable function f in  $U_{\varepsilon}(\alpha)$  and  $x_0$  sufficiently close to  $\alpha$ , the third-step method (3.6<sup>M</sup>) has eight order of convergence, i.e.,

$$e_{n+1} = (c_2^2 - c_3)c_2^2 c_4 e_n^8 - \frac{1}{2} \left[ 3c_2^5 (c_2 c_3 + 4c_4) - 2c_2^4 (3c_3^2 + 2c_5) + c_2^2 (3c_3^3 + 4c_3 c_5 + 4c_4^2) + 8c_2 c_3 c_4 (c_3 - 3c_2^2) \right] e_n^9 + O(e_n^{10}),$$

where  $e_n = x_n - \alpha$  and  $c_k$  are given in (1.4).

The corresponding computational efficiency is now much better, EFF=  $8^{1/4} \approx 1.68179$ .

# Example 3.1. Consider the equation

$$f(x) = xe^{x^2} - \sin^2 x + 3\cos x + 5 = 0,$$

with a simple zero

 $\alpha = -1.20764782713091892700941675835608409776023581894953881520592\dots$ 

In order to show the behavior of three-step methods (1.3), (3.1), (3.3), (3.6) and (3.6<sup>M</sup>) we need a multi-precision arithmetics. Starting with  $x_0 = -1$ , we use MATHEMATICA with 10000 significant digits. The errors  $e_n = x_n - \alpha$  are given in Table 3.1. Numbers in parentheses indicate decimal exponents. Besides the convergence order (r) we give also the corresponding computational efficiency (EFF).

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Table 3.1. The errors  $e_n = x_n - \alpha$ , n = 0, 1, 2, 3, 4, in three-step methods

method	(1.3)	(3.1)	(3.3)	(3.6)	$(3.6^{M})$
order	r = 8	r = 10	r = 9	r = 9	r = 8
EFF	1.41421	1.46780	1.55185	1.55185	1.68179
n = 0	2.08(-1)	2.08(-1)	2.08(-1)	2.08(-1)	2.08(-1)
n=1	-1.05(-5)	3.70(-6)	-1.19(-7)	-9.24(-8)	-2.25(-6)
n=2	-2.87(-40)	5.66(-54)	2.74(-63)	2.15(-64)	-8.57(-46)
n=3	-8.87(-317)	3.93(-532)	-5.05(-564)	-4.26(-574)	-3.77(-361)
n=4	-7.48(-2529)	1.02(-5313)	1.26(-5070)	2.204(-5161)	-5.32(-2884)

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